

## Potential theory

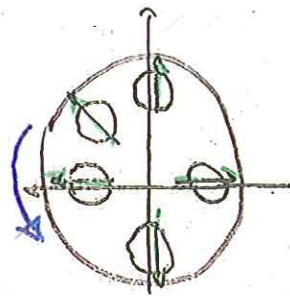
Requirements: irrotational, no rotation  
2-dimensional (planar)  
incompressible, steady

irrotational free:  $\vec{\omega} = 0$

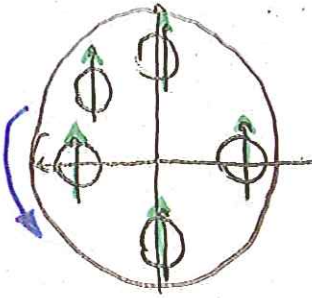
$$\vec{\omega} = \frac{1}{2} \text{rot } \vec{v} = \frac{1}{2} \nabla \times \vec{v} = \frac{1}{2} \begin{pmatrix} \omega_y - \omega_z \\ u_z - u_x \\ v_x - v_y \end{pmatrix}$$

2-dimensional flow:  $\omega_x = \omega_y = 0$

$$\rightarrow \omega_z = \frac{1}{2} (v_x - u_y) = \frac{1}{2} \left( \frac{\partial v}{\partial x} - \frac{\partial u}{\partial y} \right) = 0$$



rotational flow



irrotational free  
 $\Rightarrow$  potential flow

[Example 1.1]  $\text{rot}(\text{grad } \phi) = 0$

## Potential functions

If  $\omega_z = 0 \rightarrow$  A function  $\phi$  exists with the

property  $\vec{v} = \nabla \phi \rightarrow (v) = \left( \frac{\partial \phi}{\partial x}, \frac{\partial \phi}{\partial y} \right)$

$\Rightarrow u = \frac{\partial \phi}{\partial x}; v = \frac{\partial \phi}{\partial y}$

$\rightarrow$  Continuity (2-D, steady, incompressible)

$$\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} = \nabla \cdot \vec{v} = 0 \Rightarrow \nabla^2 \phi = \Delta \phi = 0$$

$$\Rightarrow \frac{\partial^2 \phi}{\partial x^2} + \frac{\partial^2 \phi}{\partial y^2} = 0$$

Laplace equation of the potential function

Linear differential equation

$\Rightarrow$  The principle of superposition is applicable

If  $\phi_1, \phi_2$  are solutions of the equation,

then  $C_1 \phi_1, C_2 \phi_2$  and  $C_1 \phi_1 + C_2 \phi_2$  are

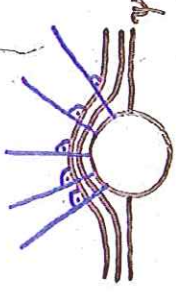
also solutions

Stream function:  $u = \frac{\partial \psi}{\partial y}; v = -\frac{\partial \psi}{\partial x}$

(fulfills the Continuity)

$$\omega = 0 \rightarrow \nabla^2 \psi = \Delta \psi = 0$$

$\phi_x = \psi_y; \phi_y = -\psi_x \rightarrow \phi$  and  $\psi$  are perpendicular



$\phi = \text{const} \rightarrow$  lines of constant potential  
 $\psi = \text{const} \rightarrow$  streamlines

$\phi$  and  $\psi$  are used to describe flow fields and flows around bodies.

The contour is provided by a special stream line  $\Rightarrow$  The velocity vector  $\vec{v}$  is parallel to the wall

But: Stokes no-slip condition cannot be fulfilled (frictionless + rotation free)

$\rightarrow$  Drag forces and viscous stresses cannot be determined

complex numbers:

$$z = x + iy = r e^{i\varphi} = r (\cos\varphi + i \sin\varphi)$$

$$x = r \cos\varphi \quad r = \sqrt{x^2 + y^2}$$

$$y = r \sin\varphi \quad \varphi = \arctan\left(\frac{y}{x}\right)$$

complex velocity  
 $w = u + iv$

conjugate complex velocity  
 $\bar{w} = u - iv$

complex potential function  
 complex stream function

$$F(z) = \int w dz = \phi(x, y) + i\psi(x, y)$$

$$\Rightarrow \text{Laplace-equation } \left( \frac{\partial^2 \phi}{\partial x^2} + \frac{\partial^2 \psi}{\partial y^2} \right) + i \left( \frac{\partial^2 \psi}{\partial x^2} + \frac{\partial^2 \phi}{\partial y^2} \right) = 0$$

$$\bar{w} = u - iv = \frac{dF}{dz}$$

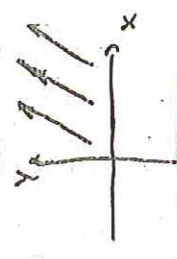
Singularities:

Parallel flow:

$$F(z) = (u_\infty - i v_\infty) z = \phi + i\psi$$

$$\phi = u_\infty x + v_\infty y \quad \psi = u_\infty y - v_\infty x$$

$$u = u_\infty \quad v = v_\infty$$



Streamlines

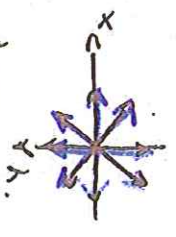
Source, sink:

$$F(z) = \frac{E}{2a} \ln z = \phi + i\psi$$

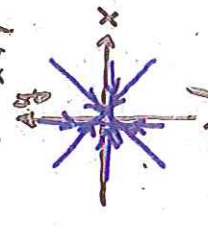
$$\phi = \frac{E}{2a} \ln r \quad \psi = \frac{E}{2a} \varphi$$

$$u = \frac{E}{2a} \frac{x}{x^2 + y^2}$$

$$v = \frac{E}{2a} \frac{y}{x^2 + y^2}$$



Source



Sink

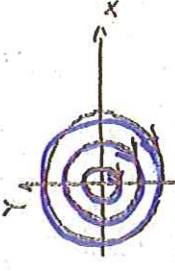


**Vortex:**

$$F(z) = \frac{\Gamma}{2\pi} i \ln z$$

$$\phi = -\frac{\Gamma}{2\pi} \arctan \frac{y}{x} \quad \psi = \frac{\Gamma}{2\pi} \ln \sqrt{x^2 + y^2}$$

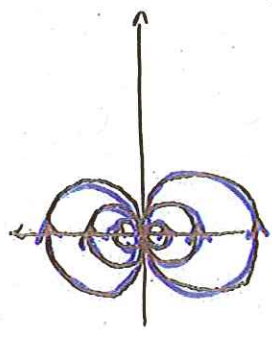
$$u = \frac{\Gamma}{2\pi} \frac{y}{x^2 + y^2} \quad v = -\frac{\Gamma}{2\pi} \frac{x}{x^2 + y^2}$$



**Dipole, Doublet**

$$F(z) = \frac{m}{z}$$

$$\phi = \frac{m}{x^2 + y^2} \quad \psi = -\frac{m}{x^2 + y^2} \quad u = m \frac{y^2 - x^2}{(x^2 + y^2)^2} \quad v = -m \frac{2xy}{x^2 + y^2}$$



Superposition of a source and a sink

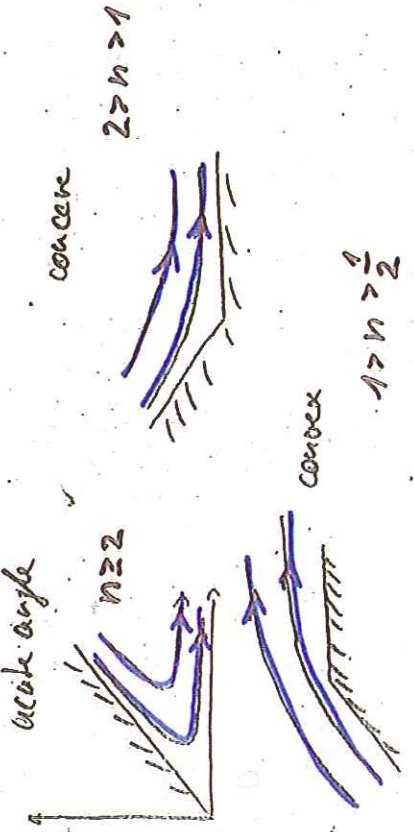
usually:  $m = \frac{M}{2\pi}$   
 $\Rightarrow F(z) = \frac{M}{2\pi z}$

**Complex Flow:**

$$F(z) = \frac{a}{n} z^n$$

$$n \in \mathbb{R} \quad a \in \mathbb{C}$$

$$\phi = \frac{a}{n} r^n \cos(n\varphi) \quad \psi = \frac{a}{n} r^n \sin(n\varphi)$$



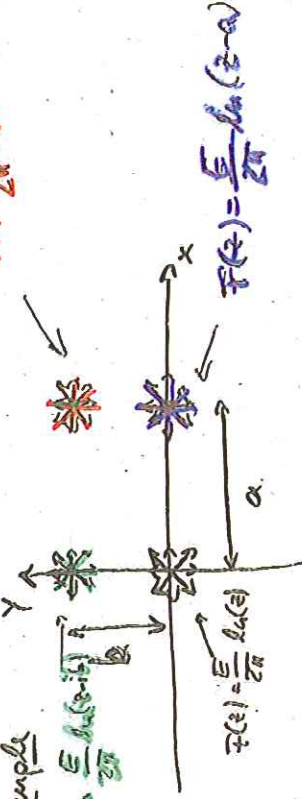
Singularities have their centers in the origin of the coordinate system

**Displacement**

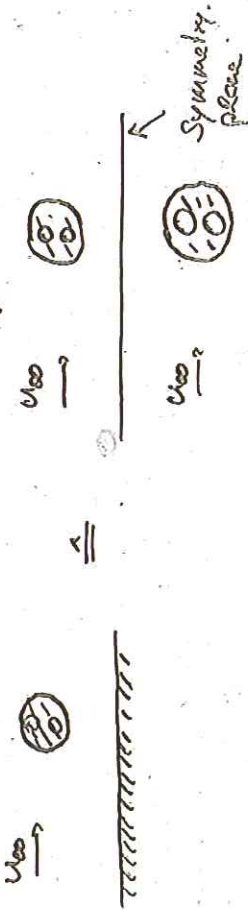
Example

$$F(z) = \frac{\Gamma}{2\pi} \ln(z-i)$$

$$F(z) = \frac{\Gamma}{2\pi} \ln(z-a-ib)$$



Simulation of wells by mirroring



Usually, the contour is presented by the stagnation streamlines

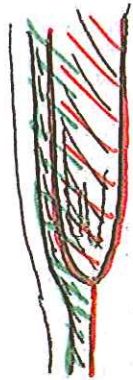
- locate the stagnation point ( $u = v = 0$ ):  $x_s, y_s$
- determine  $z$  at the stagnation point

• sketch the streamlines

$$z_k(x, y) = z_k(x_s, y_s) = \text{const.}$$

• Streamlines do not intersect

→ each streamline can represent a body contour. In this case usually  $u_{\infty} \neq 0$



• Bernoulli's equation is valid

$$p_0 = p_{\infty} + \frac{\rho}{2} \sum (u_{\infty}^2 + v_{\infty}^2) = p + \frac{\rho}{2} \sum (u^2 + v^2) = \text{const.}$$

$$\text{• Compute } q_p = \frac{p - p_{ref}}{\frac{\rho}{2} u_{ref}^2} = \frac{\frac{\rho}{2} u_{ref}^2 - \frac{\rho}{2} \sum V^2}{\frac{\rho}{2} u_{ref}^2} = 1 - \frac{\sum V^2}{u_{ref}^2}$$

14.5 A planar flow is described by the stream function  $\psi = \left(\frac{U}{L}\right)xy$ . In  $x_{ref} = 0, y_{ref} = 1$  m the pressure is  $p_{ref} = 10^5$  N/m<sup>2</sup>.

$$U = 2 \text{ m/s} \quad L = 1 \text{ m} \quad \rho = 10^3 \text{ kg/m}^3$$

a) Proof, if the flow has a potential.

Determine

b) the stagnation points, the pressure coefficient, and the lines of constant total velocity

c) the velocity and the pressure at  $x_1 = 2$  m,  $y_1 = 2$  m,

d) the coordinates of a particle at  $t = 0.5$  s if it passes at  $t = 0$  the point

$x_1, y_1$ ,

e) the pressure difference between these two points.

f) Sketch the streamlines.

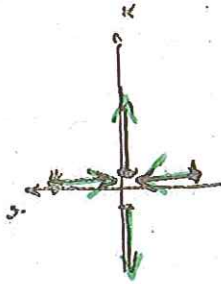
Sketch of the flow field

- Stagnation points - stagnation streamline
- asymptotic streamlines for  $x, y \rightarrow \infty$ ;  $x, y \rightarrow 0$
- flow direction

Stagnation points:  $\vec{v} = \vec{0}$ ;  $u = v = 0$

$$u = \frac{U}{L} x; v = -\frac{U}{L} y \rightarrow (x_s, y_s) = (0, 0)$$

Additionally:  $u = 0$  on the  $y$ -axis  
 $v = 0$  on the  $x$ -axis



Stream Lines:  $\zeta = \text{const}$

$$\zeta = \frac{U}{L} xy = \text{const}$$

$$\rightarrow y = \frac{L}{U} \text{const} \cdot \frac{1}{x} = \frac{C}{x} \quad \text{for } x \neq 0$$

$$x = \frac{L}{U} \text{const} \cdot \frac{1}{y} = \frac{C}{y} \quad \text{for } y \neq 0$$

$\rightarrow$  Hyperbola

Stagnation point stream line

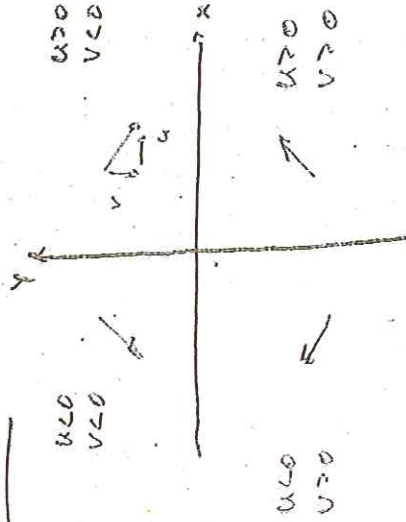
$$u_{sp} = \frac{U}{L} x_{sp} y_{sp} = 0 \quad \left| \text{dependent on the problem} \right.$$

$$\zeta = 0 \rightarrow x = 0 \quad \text{or} \quad y = 0$$

$\rightarrow$   $x$ -axis and  $y$ -axis are stagnation stream lines

$$u = \frac{U}{L} x; v = -\frac{U}{L} y$$

flow direction





a) Given: Streamfunction  $\psi = \frac{u}{L} xy$

$\phi$  exists, if  $\vec{\omega} = 0$

planar flow  $\rightarrow$  2-dimensional  $\rightarrow \omega_x = \omega_y = 0$

$$\omega_z = \frac{1}{2} \left( \frac{\partial v}{\partial x} - \frac{\partial u}{\partial y} \right)$$

$$u = \frac{\partial \psi}{\partial y} = \frac{u}{L} x \rightarrow \frac{\partial u}{\partial y} = 0 \quad \omega_z = 0$$

$$v = -\frac{\partial \psi}{\partial x} = -\frac{u}{L} y \rightarrow \frac{\partial v}{\partial x} = 0 \quad \rightarrow \text{the flow is irrotational and the potential exists}$$

$\phi$  exists  $\rightarrow$  Computation of  $\phi$

$$1) u = \frac{\partial \phi}{\partial x} \rightarrow \phi = \int u dx + f_1(y) + C_1$$

$$2) v = \frac{\partial \phi}{\partial y} \rightarrow \phi = \int v dy + f_2(x) + C_2$$

$$1) \phi(x, y) = \int \frac{u}{L} x dx + f_1(y) + C_1$$

$$2) \phi(x, y) = \int -\frac{u}{L} y dy + f_2(x) + C_2$$

$$1) \phi(x, y) = \frac{u}{L} \frac{x^2}{2} + f_1(y) + C_1$$

$$2) \phi(x, y) = -\frac{u}{L} \frac{y^2}{2} + f_2(x) + C_2$$

Comparison of 1) and 2)

$$\frac{u}{L} \frac{x^2}{2} + f_1(y) + C_1 = -\frac{u}{L} \frac{y^2}{2} + f_2(x) + C_2$$

$$f_1(y) = -\frac{u}{L} \frac{y^2}{2} \quad ; \quad f_2(x) = \frac{u}{L} \frac{x^2}{2} \quad ; \quad C_1 = C_2 = C$$

$$\rightarrow \phi = \frac{u}{2L} (x^2 - y^2) + C$$

Complex potential  $\bar{z}(z)$

$$\bar{z}(z) = \bar{z}(x+iy) = \phi(x, y) + i\bar{\psi}(x, y)$$

$$= \frac{u}{2L} (x^2 - y^2) + i \frac{u}{L} xy$$

$$= \frac{u}{2L} (x^2 + 2ixy - y^2)$$

$$= \frac{u}{2L} z^2$$

pressure coefficient

$$C_p = \frac{p - p_{\infty}}{\frac{\rho}{2} V_{\infty}^2}$$

$$\frac{p - p_{\infty}}{p_{\infty}} = 1 - \frac{V^2}{V_{\infty}^2}$$

$$\frac{u^2 + v^2}{u_{\infty}^2 + v_{\infty}^2}$$

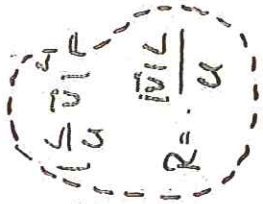
$$u = \frac{U}{L} x \quad v = -\frac{U}{L} y$$

$$C_p = 1 - \frac{x^2 + y^2}{x_{\infty}^2 + y_{\infty}^2}$$

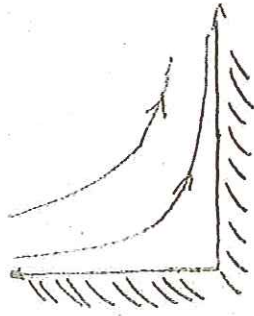
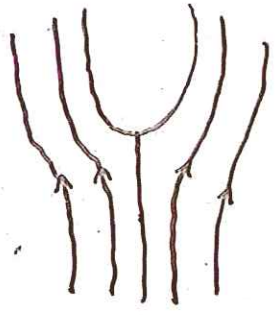
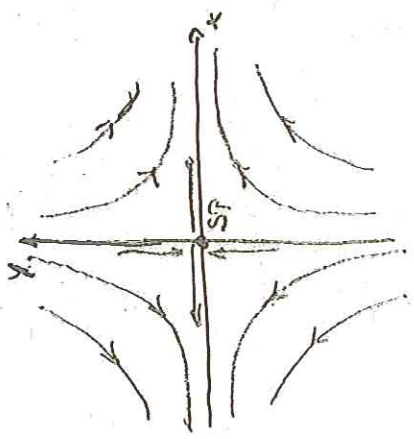
lines of constant velocity

$$|\vec{V}| = \sqrt{u^2 + v^2} = \text{const}$$

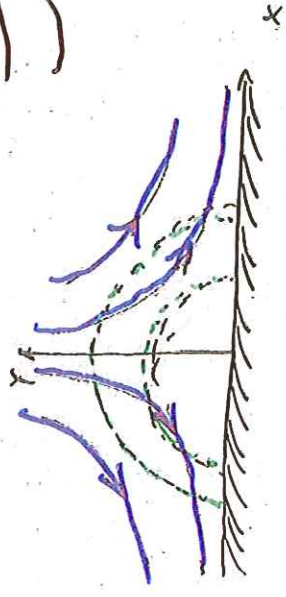
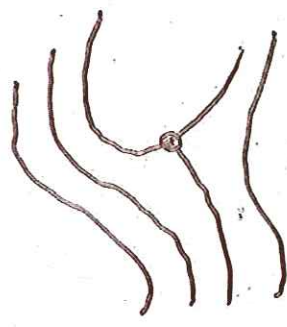
$$= \sqrt{\left(\frac{U}{L} x\right)^2 + \left(-\frac{U}{L} y\right)^2} = \left(\frac{U}{L}\right) \sqrt{x^2 + y^2}$$



circle with the radius



90° - Corner flow



flow stagnation point flow

14.4 The complex stream function is given

$$F(z) = \frac{2u_\infty}{3\sqrt{L}} z^{\frac{3}{2}} + \frac{E}{2\pi} \ln(z)$$

Given:  $L, u_\infty$

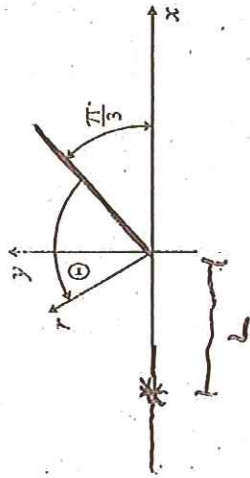
Determine

- the potential  $\phi(r, \theta)$  and the stream function  $\psi(r, \theta)$ .
- the components of the velocity  $v_r, v_\theta$ .
- the constant  $E$  such that a stagnation point is at  $(x = -L, y = 0)$ .
- the equation that describes the contour  $r_s(\theta)$ .

Hints:

$$z = x' + iy' = re^{i\theta}$$

$$v_r = \frac{\partial \phi}{\partial r} = \frac{1}{r} \frac{\partial \psi}{\partial \theta}, \quad v_\theta = \frac{\partial \psi}{\partial r} = -\frac{\partial \phi}{\partial \theta}$$



$$e^{\frac{3}{2} \ln z} = \cos \frac{3}{2} \theta + i \sin \frac{3}{2} \theta$$

3/2

a)  $F(z) = \frac{2}{3} \frac{u_\infty}{\sqrt{L}} z^{\frac{3}{2}} + \frac{E}{2\pi} \ln z = \frac{2}{3} \frac{u_\infty}{\sqrt{L}} r^{\frac{3}{2}} e^{i \frac{3}{2} \theta} + \frac{E}{2\pi} (\ln r + i \theta)$

$$\phi(r, \theta) = \text{Re}(F(z)) = \frac{2}{3} \frac{u_\infty}{\sqrt{L}} r^{\frac{3}{2}} \cos\left(\frac{3}{2}\theta\right) + \frac{E}{2\pi} \ln r$$

$$\psi(r, \theta) = \text{Im}(F(z)) = \frac{2}{3} \frac{u_\infty}{\sqrt{L}} r^{\frac{3}{2}} \sin\left(\frac{3}{2}\theta\right) + \frac{E}{2\pi} \theta$$

b)  $v_r = \frac{\partial \phi}{\partial r} = u_\infty \sqrt{\frac{r}{L}} \cos\left(\frac{3}{2}\theta\right) + \frac{E}{2\pi r}$

$$v_\theta(r, \theta) = \frac{1}{r} \frac{\partial \phi}{\partial \theta} = -u_\infty \sqrt{\frac{r}{L}} \sin\left(\frac{3}{2}\theta\right)$$

c) stagnation point at  $(x = -L, y = 0) \rightarrow (r = L, \theta = \frac{2}{3}\pi)$

$$v_\theta = -u_\infty \sqrt{\frac{L}{L}} \sin \frac{\pi}{2} = 0$$

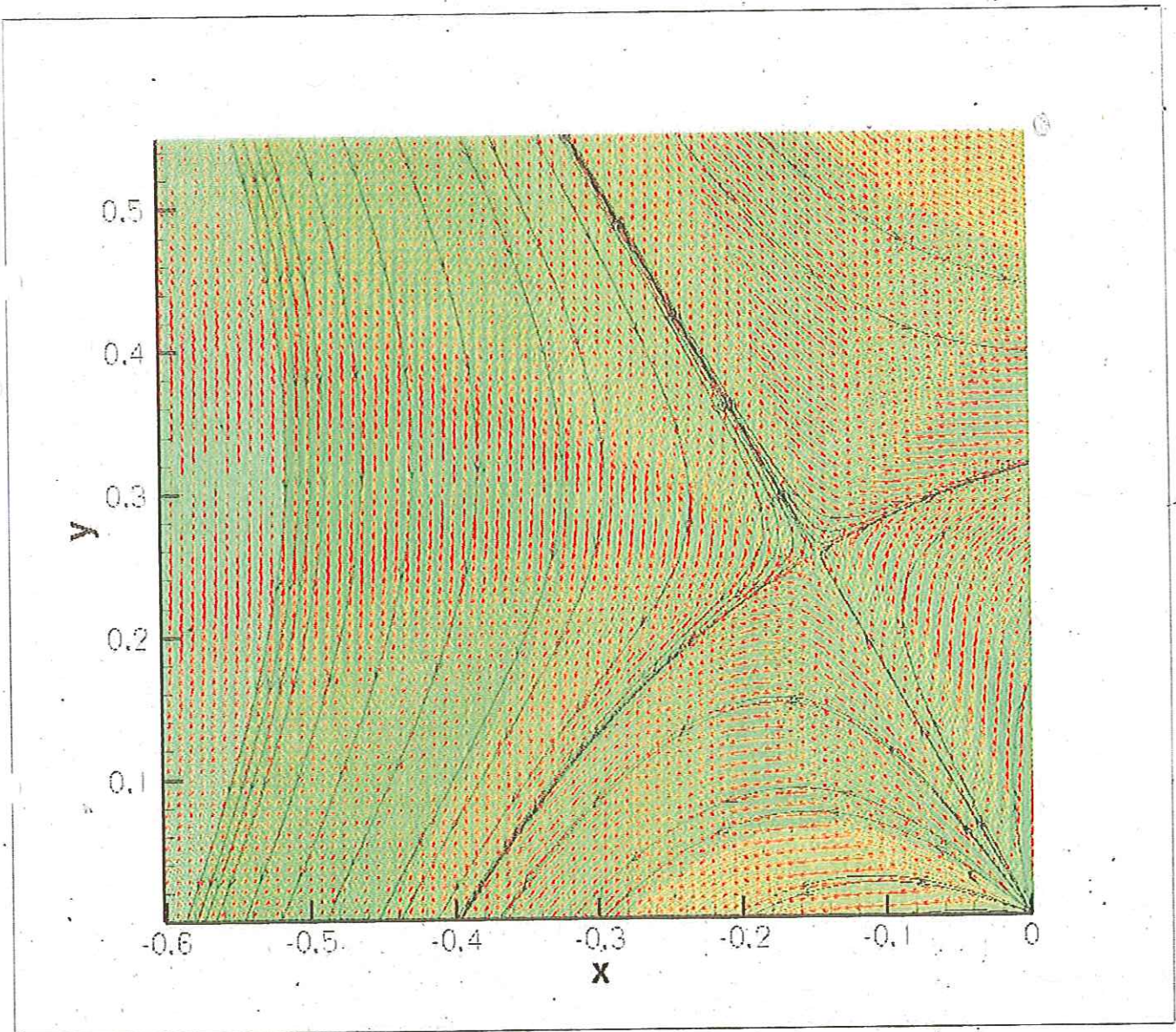
$$v_r = u_\infty \sqrt{\frac{L}{L}} \cos \frac{\pi}{2} + \frac{E}{2\pi L} = 0 \rightarrow E = 2\pi u_\infty L$$

d)  $r_{sp} = \frac{E}{2\pi} \cdot \frac{2}{3\sqrt{L}} = \frac{E}{3} = \frac{2}{3} \pi u_\infty L = 2\pi u_\infty L$

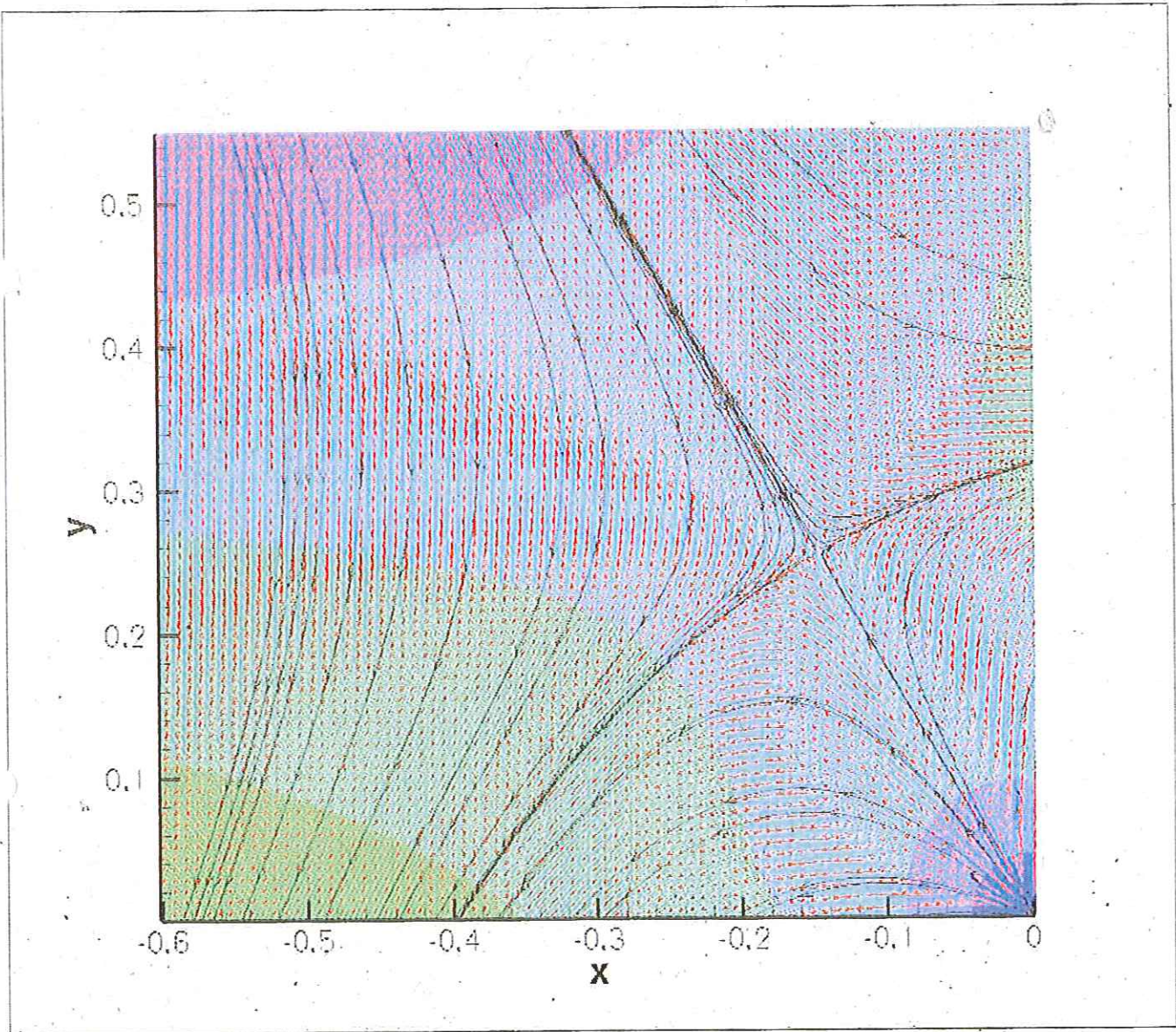
$$\frac{2}{3} \pi u_\infty L = \frac{2}{3} \frac{u_\infty}{\sqrt{L}} r^{\frac{3}{2}} \sin\left(\frac{3}{2}\theta\right) + u_\infty L \theta$$

$$\rightarrow r_s(\theta) = L \cdot \left( \frac{\pi - \frac{3}{2}\theta}{\sin \frac{3}{2}\theta} \right)^{\frac{2}{3}}$$

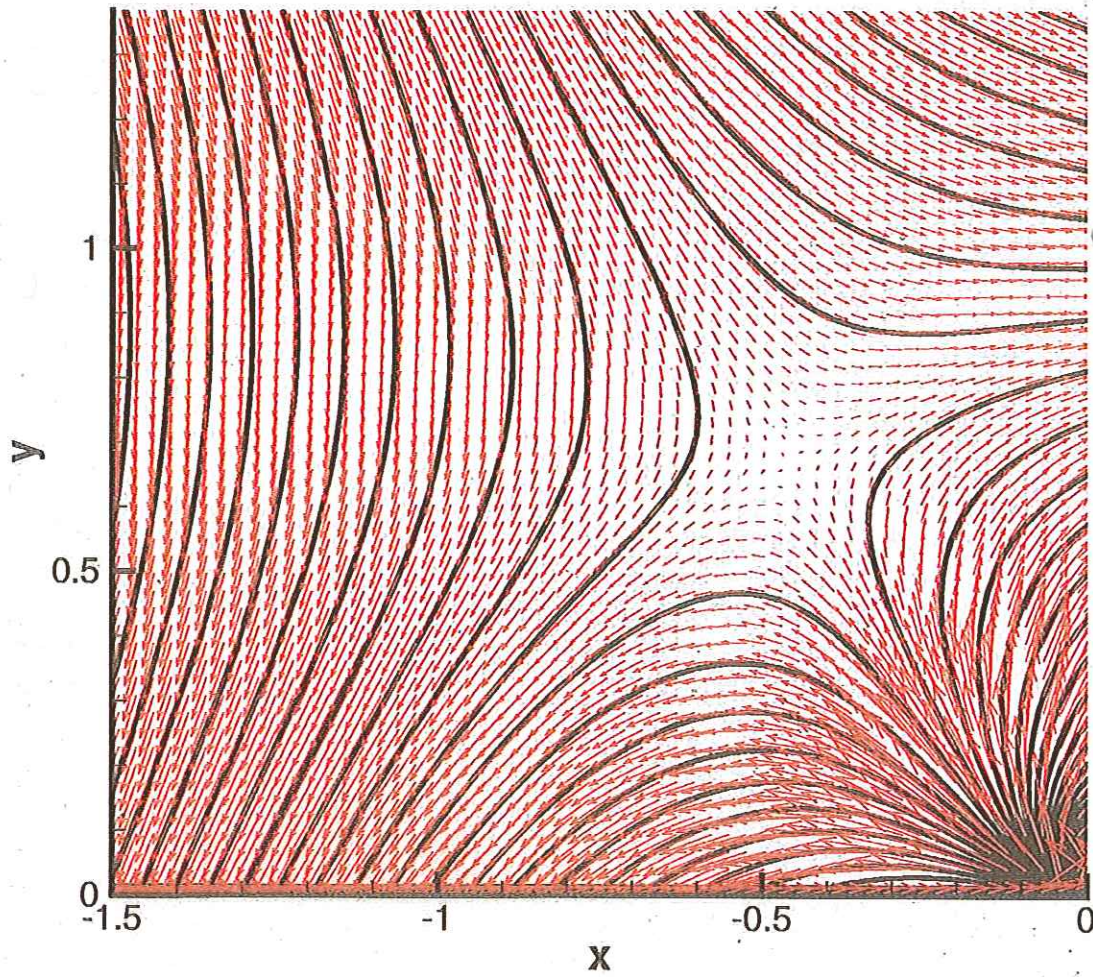






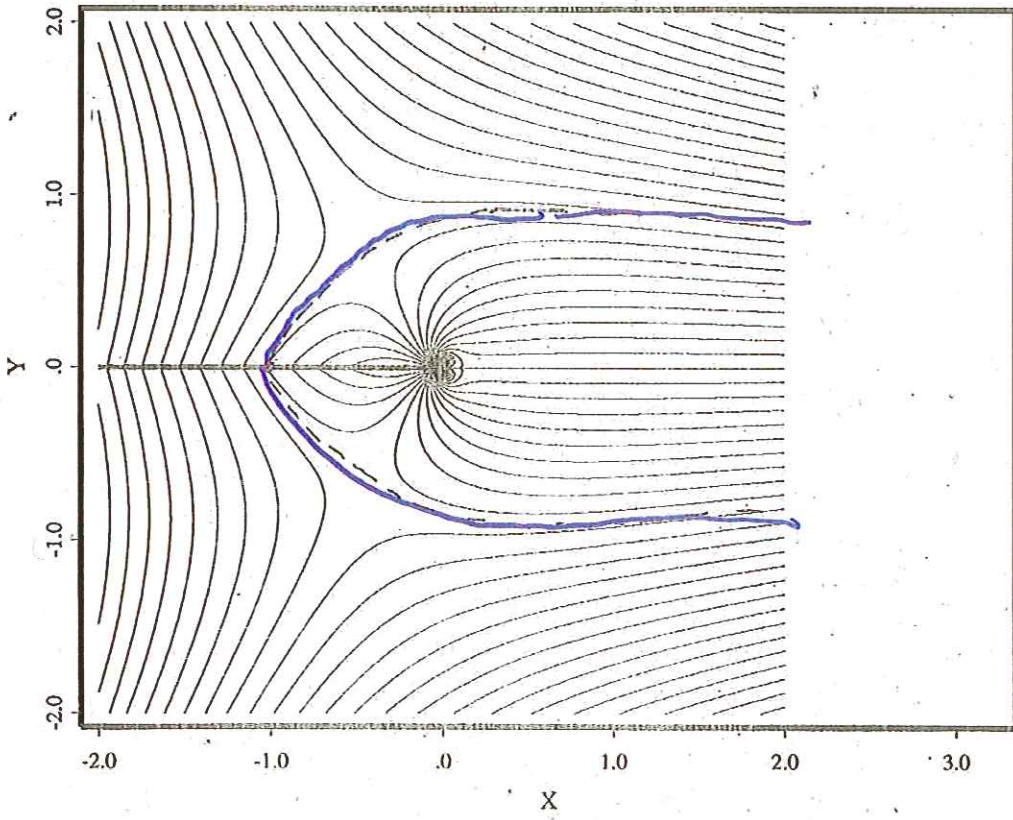




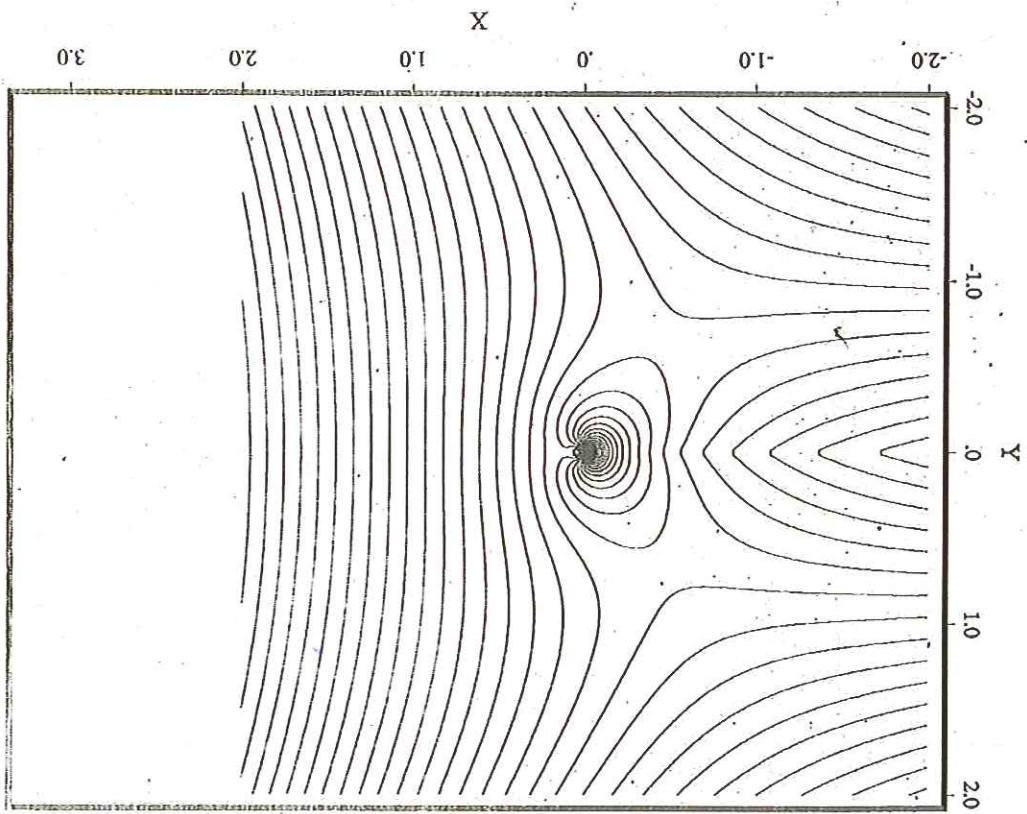




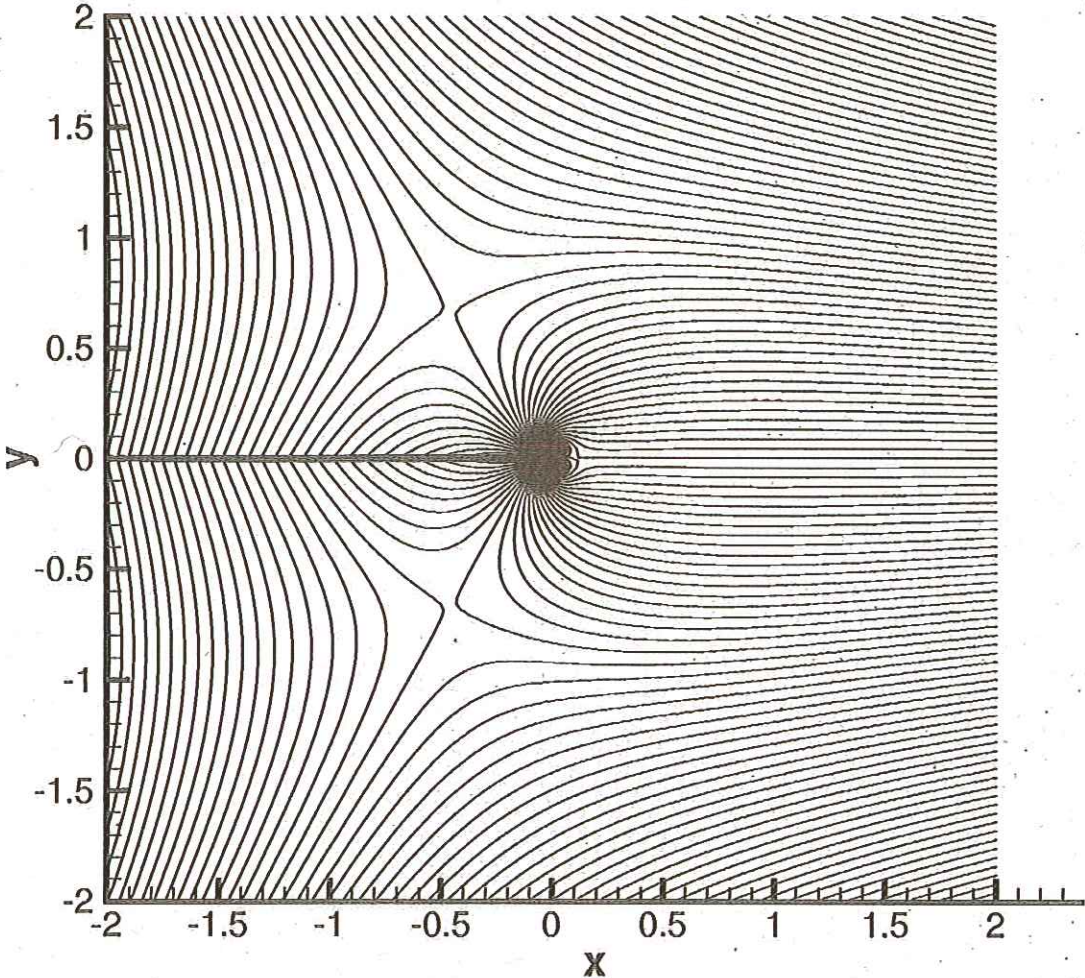
STROMFUNKTION



POTENTIAL

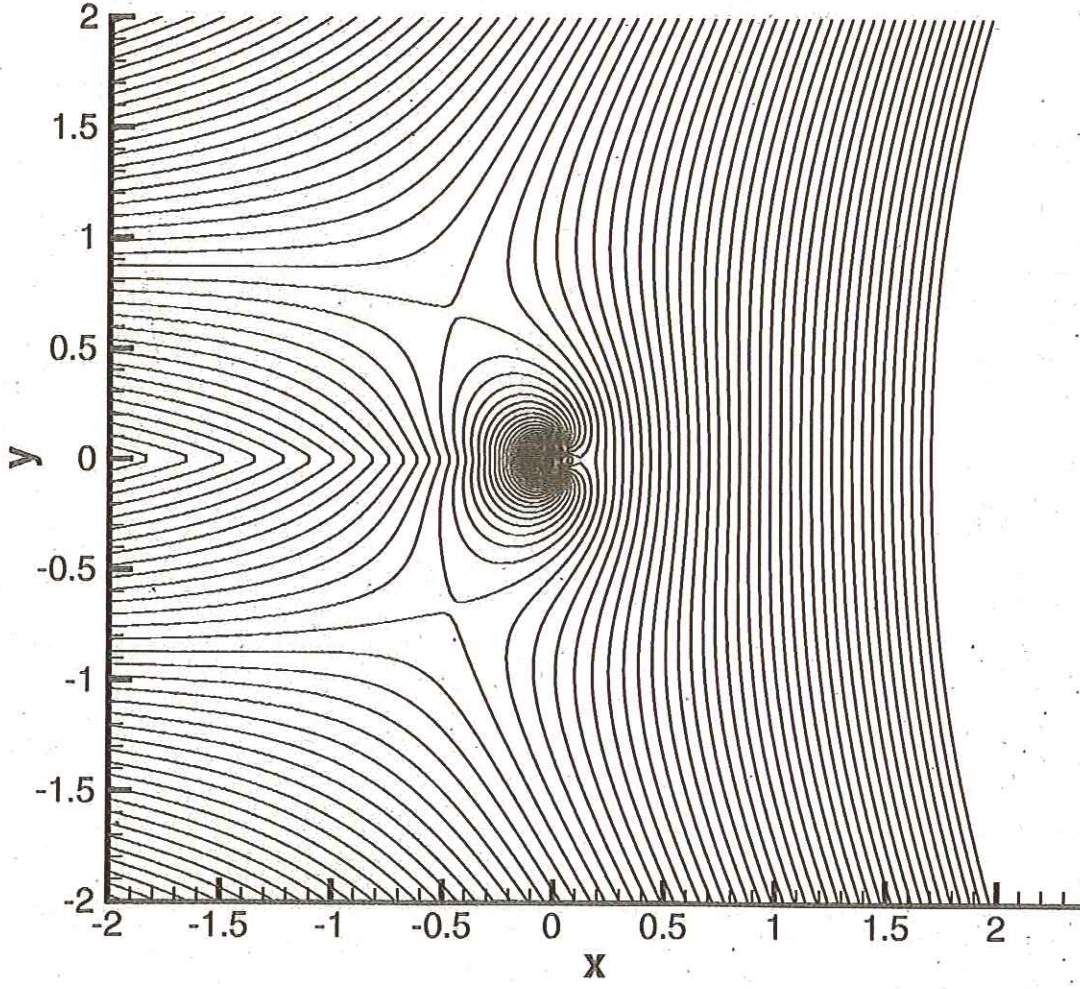


### Stromlinien





### Potentiallinien

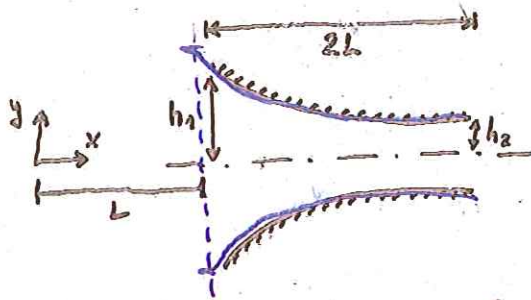




The stream function for a flow of an incompressible fluid through a planar nozzle is given:

$$\psi(x, y) = \frac{y}{h(x)} \cdot M_{\infty} \cdot L$$

given:  $M_{\infty}, L, B, h_1 = L, h_2 = \frac{1}{3}L$



a) Determine the shape  $h(x)$  of the nozzle such that the flow can be described by the potential theory.

b) Determine the velocity distribution  $u(x, y)$  and  $v(x, y)$  and the volume flux.

1) prerequisite:  $\omega \stackrel{!}{=} 0 \Rightarrow \nabla^2 \psi \stackrel{!}{=} 0$  with  $u = \frac{\partial \psi}{\partial y}$  and  $v = -\frac{\partial \psi}{\partial x}$

$$\Rightarrow \frac{\partial^2 \psi}{\partial x^2} + \frac{\partial^2 \psi}{\partial y^2} \stackrel{!}{=} 0 \quad \frac{\partial^2 \psi}{\partial y^2} = 0$$

$$\Rightarrow \frac{\partial^2 \psi}{\partial x^2} \stackrel{!}{=} 0 = \frac{\partial}{\partial x} \left( y M_{\infty} L \frac{\partial}{\partial x} \left( \frac{1}{h(x)} \right) \right)$$

$$\Rightarrow \frac{\partial}{\partial x} \left( -y M_{\infty} L \cdot \frac{h'(x)}{h^2(x)} \right) = 0 \quad \text{inner derivative}$$

$$\Rightarrow \frac{d}{dx} \left( \frac{h'(x)}{h^2(x)} \right) = 0$$

$$\Rightarrow \text{integrate twice: } -\frac{1}{h(x)} = C_1 x + C_2$$

$$\text{B.C.: } x=L, h=L \Rightarrow -\frac{1}{L} = C_1 \cdot L + C_2$$

$$x=3L, h=\frac{1}{3}L \Rightarrow -\frac{3}{L} = 3C_1 \cdot L + C_2$$

$$\left. \begin{array}{l} \\ \\ \end{array} \right\} C_2 = 0, C_1 = -\frac{1}{L^2}$$

$$\Rightarrow \boxed{h(x) = \frac{L^2}{x}}$$

b)  $u = \frac{\partial \psi}{\partial y}; v = -\frac{\partial \psi}{\partial x}$  with  $\psi = \frac{M_{\infty}}{L} x y$

$$\Rightarrow \boxed{u = M_{\infty} \frac{x}{L}} \quad \text{and} \quad \boxed{v = -M_{\infty} \frac{y}{L}}$$

$$\dot{Q} = \psi(y=h) - \psi(y=-h) \Big|_{x=L} \Rightarrow \dot{Q} = (h_1 + h_1) B \cdot M_{\infty} = \boxed{2 M_{\infty} L B}$$

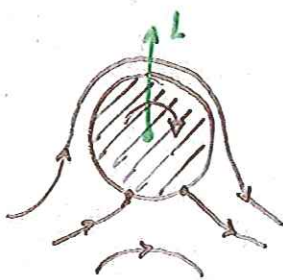
Rotating cylinders:  $\rightarrow$  Force in lateral direction  
like the lift force on an airfoil

$\Rightarrow$  potential theory can be used to describe airfoil flows

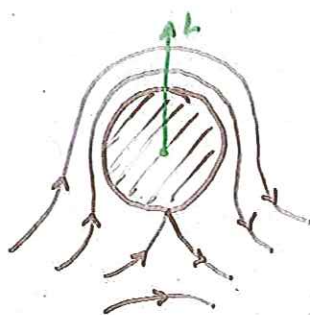
cylinders:  $F(z) = M_{\infty} z + \frac{M}{2\pi z}$

rotating cylinders:  $F(z) = M_{\infty} z + \frac{M}{2\pi z} + i \frac{\Gamma}{2\pi} \ln z$

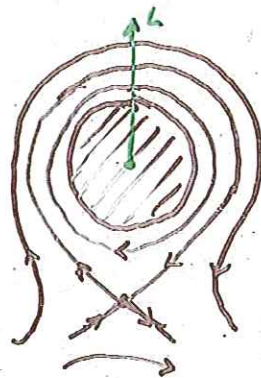
with  $M = 2\pi R^2 M_{\infty}$



$\Gamma < 4\pi M_{\infty} R$



$\Gamma = 4\pi M_{\infty} R$



$\Gamma > 4\pi M_{\infty} R$

effect of Kutta's

$\rightarrow$  the stronger  $\Gamma$ , the faster the velocities on the upper surface  $\rightarrow$  pressure drop

$\rightarrow$  pressure difference  $\Rightarrow$  lift force:  $L = \rho M_{\infty} \Gamma$

hint:  $2 \cdot \sin \varphi \cdot \cos \varphi = \sin 2\varphi$  and  $-\sin 2\varphi = \sin 2(\pi - \varphi) = \sin 2\varphi'$

$$\Rightarrow c_{pk} = \frac{\sin 2\varphi'}{\varphi'} - \left(\frac{\sin \varphi'}{\varphi'}\right)^2$$

d) isobars:  $c_p = \text{const}$

$$c_p = -\frac{h}{\pi} \cdot \frac{2x + \frac{h}{\pi}}{x^2 + y^2} \Rightarrow x^2 + y^2 = -2\frac{h}{\pi} \frac{x}{c_p} - \left(\frac{h}{\pi}\right)^2 \cdot \frac{1}{c_p} + \left(\frac{h}{\pi}\right)^2 \cdot \frac{1}{c_p^2} - \left(\frac{h}{\pi}\right) \cdot \frac{1}{c_p^2}$$

$$\Rightarrow x^2 + 2\frac{h}{\pi} \frac{x}{c_p} + \left(\frac{h}{\pi}\right)^2 \cdot \frac{1}{c_p^2} + y^2 = -\frac{h^2}{\pi^2} \cdot \frac{1}{c_p} + \frac{h^2}{\pi^2} \cdot \frac{1}{c_p^2}$$

$$\Rightarrow \left(x + \frac{h}{\pi \cdot c_p}\right)^2 + y^2 = \underbrace{(1 - c_p)}_{R^2} \left(\frac{h}{\pi c_p}\right)^2$$

$\Rightarrow$  circles around  $\left(-\frac{h}{\pi c_p}, 0\right)$  with radius  $\frac{h}{\pi c_p} \cdot \sqrt{1 - c_p}$



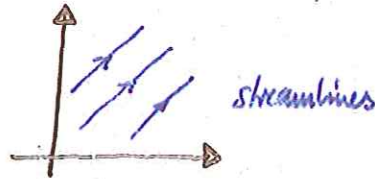
# Potential theory

$$F(z) = \int \bar{w}(z) dz = \phi(x, y) + i \psi(x, y)$$

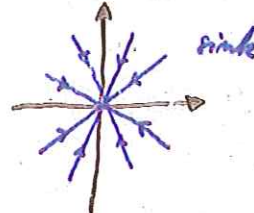
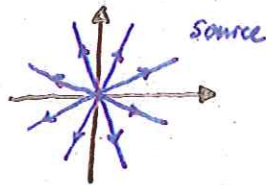
$$z = x + iy = r \cdot e^{i\varphi} \\ = r(\cos\varphi + i \sin\varphi)$$

$$\rightarrow \bar{w} = \frac{dF}{dz} = u - iv \quad \begin{array}{l} \text{potential function} \\ \text{stream function} \end{array}$$

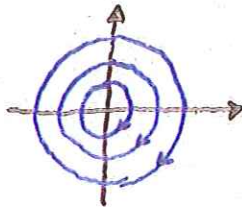
Parallel flow:  $F(z) = (u_{\infty} - iv_{\infty}) \cdot z \rightarrow \phi = u_{\infty}x + v_{\infty}y ; \psi = u_{\infty}y - v_{\infty}x$



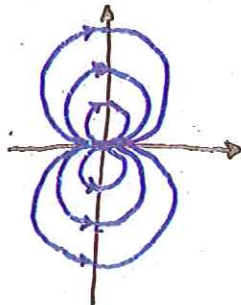
Source, sink:  $F(z) = \frac{E}{2\pi} \ln z \rightarrow \phi = \frac{E}{2\pi} \ln r ; \psi = \frac{E}{2\pi} \varphi$



Vortex:  $F(z) = \frac{\Gamma}{2\pi} i \cdot \ln z \rightarrow \phi = -\frac{\Gamma}{2\pi} \arctan \frac{y}{x} ; \psi = \frac{\Gamma}{2\pi} \ln \sqrt{x^2 + y^2} \\ = -\frac{\Gamma}{2\pi} \varphi \quad = \frac{\Gamma}{2\pi} \ln r$



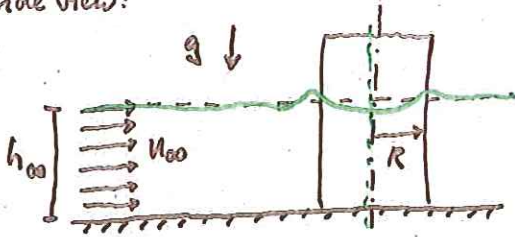
Dipole:  $F(z) = \frac{M}{2\pi z} \rightarrow \phi = \frac{M}{2\pi} \frac{x}{x^2 + y^2} ; \psi = \frac{-M}{2\pi} \frac{y}{x^2 + y^2}$



$\rightarrow$  superposition of sink and source

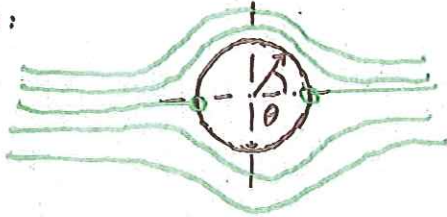
Water is flowing against a bridge pier with the velocity  $u_{\infty}$ . The pier has a circular cross-section. Far away from the pier the water level is  $h_{\infty}$ .

side view:



$$\begin{aligned} u_{\infty} &= 1 \text{ m/s} \\ h_{\infty} &= 6 \text{ m} \\ R &= 2 \text{ m} \\ \rho &= 10^3 \text{ kg/m}^3 \\ g &= 10 \text{ m/s}^2 \end{aligned}$$

top view:



- Determine the water level at the pier as a function of  $\theta$ .
- Calculate the water level in the stagnation points.
- Calculate the minimum water level.

circular cylinders: Dipole + parallel flow

$$\Rightarrow F(z) = u_{\infty} z + \frac{M}{2\pi z} \quad \text{with } z = r(\cos\theta + i\sin\theta)$$

$$\begin{aligned} \Rightarrow \phi &= \left(u_{\infty} + \frac{M}{2\pi r^2}\right) r \cos\theta \\ \psi &= \left(u_{\infty} - \frac{M}{2\pi r^2}\right) r \sin\theta \end{aligned}$$

$$\begin{aligned} \text{with } \frac{1}{z} &= \frac{1}{r(\cos\theta + i\sin\theta)} \cdot \frac{r(\cos\theta - i\sin\theta)}{r(\cos\theta - i\sin\theta)} \\ &= \frac{r(\cos\theta - i\sin\theta)}{r^2} \end{aligned}$$

$$\Rightarrow v_r = \frac{\partial\phi}{\partial r} = \left(u_{\infty} - \frac{M}{2\pi r^2}\right) \cos\theta$$

$$\Rightarrow v_{\theta} = \frac{1}{r} \frac{\partial\phi}{\partial\theta} = -\left(u_{\infty} + \frac{M}{2\pi r^2}\right) \sin\theta$$

symmetric flow: stagnation points at  $\theta=0$  and  $\theta=\pi$  and  $r=R$ :

$$\Rightarrow v_r = v_{\theta} = 0 \Rightarrow u_{\infty} - \frac{M}{2\pi R^2} = 0 \Rightarrow M = u_{\infty} \cdot 2\pi \cdot R^2$$

$$\Rightarrow v_{\theta} = -\left(u_{\infty} + \frac{R^2 u_{\infty}}{r^2}\right) \sin\theta$$

On the surface of the cylinder:  $v_r = 0$  and  $v_{\theta} = -2u_{\infty} \sin\theta$

Bernoulli:  $p_a + \rho g h_{\infty} + \frac{\rho}{2} u_{\infty}^2 = p_a + \rho g h(\theta) + \frac{\rho}{2} \vec{v}^2$

$$\rightarrow \vec{v}^2 = v_{\theta}^2 = 4u_{\infty}^2 \sin^2\theta$$

$$\Rightarrow h(\theta) = h_{\infty} + \frac{u_{\infty}^2}{2g} (1 - 4\sin^2\theta)$$

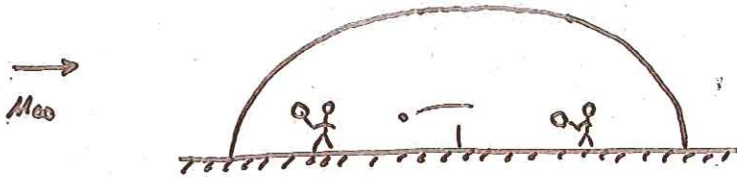
$$\begin{aligned} h(0, \pi) &= \\ &= h_{\infty} + \frac{u_{\infty}^2}{2g} = 6.05 \text{ m} \end{aligned}$$

$$\theta_{\min} = \frac{\pi}{2}, \frac{3\pi}{2}$$

$$\begin{aligned} h\left(\frac{\pi}{2}, \frac{3\pi}{2}\right) &= \\ &= h_{\infty} - \frac{3u_{\infty}^2}{2g} = 5.85 \text{ m} \end{aligned}$$



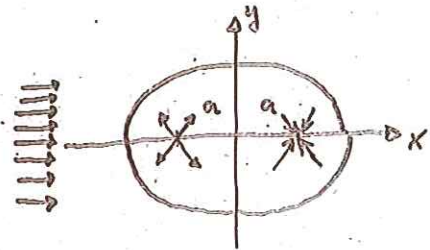
The flow around a tennis sports hall should be described by potential theory.



- What potential function should be used to describe the problem?
- Determine the velocities  $u(x,y)$  and  $v(x,y)$  using the complex conjugate velocity  $\bar{w}$ .
- Determine the stagnation points and the contour streamline.
- Sketch the streamlines.
- What 2 simplifications are used in potential theory according to real flows?

a) Superposition of parallel flow, source and sink:

$$F(z) = M_{\infty} z + \frac{E}{2\pi} \ln(z+a) - \frac{E}{2\pi} \ln(z-a)$$



$$b) \bar{w} = \frac{dF}{dz} = u - iv = M_{\infty} + \frac{E}{2\pi} \left( \frac{1}{z+a} - \frac{1}{z-a} \right)$$

$$= M_{\infty} + \frac{E}{2\pi} \left( \frac{x+a-iy}{(x+a)^2+y^2} - \frac{x-a-iy}{(x-a)^2+y^2} \right)$$

$$\Rightarrow u = M_{\infty} + \frac{E}{2\pi} \left( \frac{x+a}{(x+a)^2+y^2} - \frac{x-a}{(x-a)^2+y^2} \right)$$

$$v = \frac{E}{2\pi} \left( \frac{y}{(x+a)^2+y^2} - \frac{y}{(x-a)^2+y^2} \right)$$

c) stagnation point is part of the contour streamline

s.p.  $u=v=0 \rightarrow$  symmetry:  $y_s = 0 \rightarrow v(x_s, y_s) = 0$

$$u=0 = M_{\infty} + \frac{E}{2\pi} \left( \frac{x_s+a}{(x_s+a)^2+y_s^2} - \frac{x_s-a}{(x_s-a)^2+y_s^2} \right) = M_{\infty} + \frac{E}{2\pi} \left( \frac{1}{x_s+a} - \frac{1}{x_s-a} \right)$$

$$= M_{\infty} - \frac{E}{2\pi} \left( \frac{2a}{x_s^2-a^2} \right)$$

$$\Rightarrow x_s = \pm \sqrt{\frac{\alpha E}{\pi M_{\infty}} + a^2}$$



contour stream line:  $\psi(x, y) = \psi_k = \text{Im}(F(z))$

$$F(z) = \phi + i\psi$$

$$= M_{\infty} (x + iy) + \frac{\Gamma}{2\pi} \left[ \ln r_1 e^{i\varphi_1} - \ln r_2 e^{i\varphi_2} \right]$$

$$\text{with: } r_1 = \sqrt{(x+a)^2 + y^2} \quad ; \quad r_2 = \sqrt{(x-a)^2 + y^2}$$

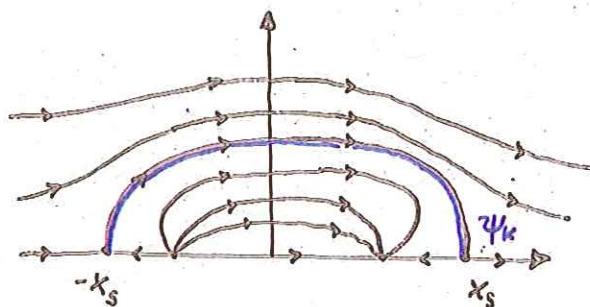
$$\varphi_1 = \arctan\left(\frac{y}{x+a}\right) \quad ; \quad \varphi_2 = \arctan\left(\frac{y}{x-a}\right)$$

$$\psi(x, y) = M_{\infty} y + \frac{\Gamma}{2\pi} \cdot \left( \arctan \frac{y}{x+a} - \arctan \frac{y}{x-a} \right)$$

$$\psi_k = \psi_{sp} \Rightarrow \psi_{sp}(x_s, y_s) = 0 \quad \text{since } y_s = 0$$

$$\Rightarrow 0 = M_{\infty} y + \frac{\Gamma}{2\pi} \cdot \left( \arctan \frac{y}{x+a} - \arctan \frac{y}{x-a} \right)$$

d) sketch:



- e) simplifications:
- contour is equal to streamline  $\rightarrow$  no-slip condition not fulfilled
  - rotation free flow is assumed